



# Extended and broad Ly $\alpha$ emission around a BAL quasar at $z \sim 5$

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Accepted 2018 February 5. Received 2018 January 26; in original form 2018 January 9

## ABSTRACT

In this work we report deep MUSE observations of a broad absorption line (BAL) quasar at  $z \sim 5$ , revealing a Ly  $\alpha$  nebula with a maximum projected linear size of  $\sim 60$  kpc around the quasar (down to our  $2\sigma$  SB limit per layer of  $\sim 9 \times 10^{-19}$  erg s $^{-1}$  cm $^{-2}$  arcsec $^{-2}$  for a 1 arcsec $^2$  aperture). After correcting for the cosmological surface brightness dimming, we find that our nebula, at  $z \sim 5$ , has an intrinsically less extended Ly  $\alpha$  emission than nebulae at lower redshift. However, such a discrepancy is greatly reduced when referring to comoving distances, which take into account the cosmological growth of dark matter (DM) haloes, suggesting a positive correlation between the size of Ly  $\alpha$  nebulae and the sizes of DM haloes/structures around quasars. Differently from the typical nebulae around radio-quiet non-BAL quasars, in the inner regions ( $\sim 10$  kpc) of the circumgalactic medium of our source, the velocity dispersion of the Ly  $\alpha$  emission is very high (FWHM  $> 1000$  km s $^{-1}$ ), suggesting that in our case we may be probing outflowing material associated with the quasar.

**Key words:** galaxies: haloes – intergalactic medium – quasars: emission lines – quasars: general – cosmology: observations.

## 1 INTRODUCTION

The circumgalactic medium (CGM) is tightly linked to the evolutionary phases of galaxies, both by hosting the gas reservoir necessary to fuel star formation and also being subject to galactic feedback and recycling. Thus, observing the nature of the CGM and studying its interplay with the host galaxy, especially at high redshift, helps us to understand how galaxies evolve and how their evolution can influence (and be influenced by) the surrounding environment.

A common observational technique used to characterize the cold gas phase of the CGM has been the analysis of its absorption signatures along background sightlines (Hennawi et al. 2006; Prochaska & Hennawi 2009; Bouché et al. 2013, 2016; Farina et al. 2013; Hennawi & Prochaska 2013; Prochaska et al. 2017) which has provided important statistical constraints on several properties of the CGM (Adelberger et al. 2005; Steidel et al. 2010; Prochaska et al. 2013; Lau, Prochaska & Hennawi 2016, 2017). However, this method is limited by the sparseness of bright background sources and by its one-dimensional nature, which precludes from capturing the spatial distribution of the CGM.

An alternative approach is to map the CGM through direct imaging of the Ly  $\alpha$  line. Theoretical models suggest that three main mechanisms should be able to generate circumgalactic Ly  $\alpha$  emission: cooling radiation of gravitationally heated gas (e.g. Haiman, Spaans & Quataert 2000; Yang et al. 2006; Dijkstra & Loeb 2009), ultraviolet (UV) photons produced through shock mechanisms (Taniguchi & Shioya 2000; Mori, Umemura & Ferrara 2004), and recombination radiation following photoionization (often referred as *fluorescence*) powered by UV sources (Cantalupo et al. 2005; Geach et al. 2009; Kollmeier et al. 2010). While the fluorescent signal powered by the diffuse metagalactic UV background (Hogan & Weymann 1987; Binette et al. 1993; Gould & Weinberg 1996; Haardt & Madau 1996), with an expected surface brightness (SB) of  $SB_{Ly\alpha} \sim 10^{-20}$  erg s $^{-1}$  cm $^{-2}$  arcsec $^{-2}$  (Cantalupo et al. 2005; Rauch et al. 2008), is still out of reach for current optical instrumentation (but see Gallego et al. 2018), Ly  $\alpha$  fluorescence is predicted to be boosted up into the detectable regime in the vicinity of bright ionizing sources, as luminous quasars (Rees 1988; Haiman & Rees 2001; Alam & Miralda-Escudé 2002; Cantalupo et al. 2005). This theoretical prediction has been confirmed by a number of surveys targeting the fluorescent Ly  $\alpha$  emission around luminous and radio-quiet quasars, using narrow-band (NB) filters on 8-metre class optical telescopes (e.g. Cantalupo, Lilly & Haehnelt 2012; Cantalupo et al.

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2014; Martin et al. 2014; Arrigoni Battaia et al. 2016) and spectroscopic observations (e.g. Christensen et al. 2006; North et al. 2012; Herenz et al. 2015). However, these surveys have revealed giant Ly  $\alpha$  nebulae, spanning distances from the quasar larger than 100 physical kpc (pkpc), only in less than 10 per cent of the targets (e.g. Cantalupo et al. 2014; Hennawi et al. 2015; maximum projected linear sizes  $>300$  pkpc), and emission on smaller scales ( $R \lesssim 50$ – $60$  pkpc) have been detected only in about 50 per cent of the cases. This 50 per cent detection rate is likely due to a combination of limits of the observational techniques, as for instance NB filter losses, spectroscopic slit losses, point spread function (PSF) losses and, most importantly, dilution of the signal of the Ly  $\alpha$  line into the continuum flux (both background and from the host galaxy) encompassed by the width of the filter (see also Borisova et al. 2016 for a discussion).

Only recently, extraordinary advances have been made using the new Multi Unit Spectroscopic Explorer (MUSE; Bacon et al. 2010), a wide integral field spectrograph mounted on the ESO/VLT. The ability of MUSE in overcoming the technical limitations of previous observational techniques, combined with its large field of view (FOV; 1 arcmin  $\times$  1 arcmin) have boosted the detection rate up to 100 per cent, providing ubiquitous detections of Ly  $\alpha$  nebulae with maximum projected linear sizes of  $\sim 100$  pkpc around bright quasars at  $3 < z < 4$  in only 1 h observation per source (see Borisova et al. 2016 and Arrigoni-Battaia et al., in preparation).

In this study we exploit MUSE to extend the analysis of Ly  $\alpha$  emission surrounding quasars towards higher redshifts and examine the nearly unexplored (within this context) class of broad absorption line (BAL) quasars (another case has been recently reported by Arrigoni Battaia et al. 2018, i.e. SDSS J1020+1040 at  $z = 3.167$ ). In particular, we report deep ( $\sim 4$  h) MUSE observations of J1605–0112, a BAL quasar at  $z \sim 4.92$ . BALs are typically characterized by deep and blueshifted (by thousands  $\text{km s}^{-1}$ ) broad absorption troughs associated with UV resonant lines (typically C IV and Si IV) tracing very fast outflows (e.g. Weymann et al. 1991; Maiolino et al. 2004; Dunn et al. 2010), and they are thought to mark a specific phase of galaxy evolution in which the quasar wind is particularly powerful.

We present our results in the following order. In Section 2 we describe our target, the basic steps of the data reduction processes, and the analysis performed to search for extended Ly  $\alpha$  emission. In Section 3 we report the results of this work and in Section 4 a discussion of their implications. Conclusions are summarized in Section 5. Throughout the paper, we assume a Lambda cold dark matter ( $\Lambda$ CDM) cosmology with  $\Omega_m = 0.3$ ,  $\Omega_\Lambda = 0.7$ , and  $h = 70 \text{ km s}^{-1}$ . 1 arcsec at  $z \sim 4.92$  corresponds to  $\sim 6.33$  pkpc.

## 2 OBSERVATIONS AND DATA PROCESSING

J1605–0112, a radio-quiet quasar at  $z = 4.92$ , is one of the most extreme low ionization lines-BAL quasar in the Universe, both in terms of luminosity and in terms of depth and velocity of the absorption features (see Table 1). Indeed, the C IV 1549 Å absorption troughs are blueshifted by more than  $30\,000 \text{ km s}^{-1}$ , tracing very high velocity outflowing gas, and have a balnicity index<sup>1</sup> of  $9300 \text{ km s}^{-1}$  (see Maiolino et al. 2004). A full 1D spectrum obtained with MUSE and a discussion on the BAL features are reported in the Appendix A.

<sup>1</sup> The balnicity index is similar to the Equivalent Width expressed in  $\text{km s}^{-1}$ , as defined in Weymann et al. (1991).

**Table 1.** Summary of the properties of J1605–0112.

RA J2000	Dec. J2000	Redshift	$\log(\lambda L_\lambda)^a$	BI(C IV) <sup>b</sup> ( $\text{km s}^{-1}$ )
16:05:01.2	−01:12:20.6	4.92	46.46	$9300 \pm 200$

<sup>a</sup> Rest-frame luminosity  $\lambda L_\lambda$  at 1450 Å (Maiolino et al. 2003).

<sup>b</sup> Balnicity index of the C IV absorption (Maiolino et al. 2004).

MUSE observations of J1605–0112 have been acquired as part of the ESO programme ID 095.A-0875 (PI: R. Maiolino) between 2015 May and August at the UT4 ESO/VLT. The collection of data was split into five observing blocks, each of them composed by four exposures of 850 s each, with the exposures rotated with respect to each other by  $90^\circ$ . The total integration time on source was  $\sim 4$  h. Observations were carried out under good seeing conditions (FWHM  $\sim 0.6$  arcsec) and airmass  $< 1.5$ . We achieve an SB limit of  $\sim 5 \times 10^{-19} \text{ erg s}^{-1} \text{ cm}^{-2} \text{ arcsec}^{-2}$  ( $1\sigma$ ) for a  $1 \text{ arcsec}^2$  aperture in a single wavelength layer (i.e.  $1.25 \text{ Å}$ ). The SB limit on a pseudo-NB image of  $30 \text{ Å}$  (a common size for filters used in the past NB surveys), obtained averaging 24 wavelength layers around the location of the Ly  $\alpha$ , is  $\sim 3 \times 10^{-18} \text{ erg s}^{-1} \text{ cm}^{-2} \text{ arcsec}^{-2}$  ( $1\sigma$ ).

### 2.1 Data reduction

Individual exposures have been processed with basic data reduction techniques using the standard ESO–MUSE pipeline v1.6 (Weilbacher et al. 2012, 2014). For each of them we performed bias subtraction, flat fielding, twilight and illumination correction, and wavelength calibration, using the MUSE\_BIAS, MUSE\_FLAT, and MUSE\_WAVECAL pipeline recipes. Following Borisova et al. (2016) we did not perform the sky subtraction using the pipeline receipt but we postponed this task to a later stage, as discussed below. We then registered the individual exposures using the position of point sources in the field; in this way we ensure an accurate relative astrometry, as shifts of a few tenths of arcseconds can occur owing to the spatial shifts introduced by the derotator wobble between exposures (Bacon et al. 2015). The next steps of the data-reduction process were performed with CubeFix and CubeSharp, custom tools for improved flat fielding correction and sky subtraction respectively, both parts of the CUBEXTRACTOR software package (Cantalupo, in preparation). These tools, extensively described in Borisova et al. (2016) and Fumagalli et al. (2017), have been specifically developed to improve data quality for the detection of faint and diffuse emission in MUSE data cubes. As last step, we combined all the corrected and sky-subtracted data cubes of individual exposures using an average  $3\sigma$ -clipping algorithm. At this point we created a white-light image (obtained by collapsing the data cube along the wavelength direction) of the combined data cube in order to identify the continuum sources. We used the positions and spectra of the detected continuum sources to perform another iteration of flat fielding correction and sky subtraction on individual exposures, before combining them again. One further iteration is enough to improve substantially the removal of self-calibration effects and data quality. The final data cube samples the instrument FOV of  $1 \times 1 \text{ arcmin}^2$  in pixels of size  $0.2 \text{ arcsec}$  and, for each pixel, contains a spectrum covering the wavelength range  $4750$ – $9350 \text{ Å}$  in spectral bins of  $1.25 \text{ Å}$ .

In Appendix B we show the whole  $1 \times 1 \text{ arcmin}^2$  FOV of the final data cube. To check the accuracy of flux calibration we compared our final product with photometric data from the Sloan Digital Sky

Survey (SDSS; Eisenstein et al. 2011) without finding any flux deviation within the errors for the few sources in common.

## 2.2 Data analysis

### 2.2.1 PSF and continuum subtraction

Revealing extended Ly $\alpha$  structures around a quasar requires a proper subtraction of the PSF. Generally, accurate PSF removal is a substantial issue in the case of NB imaging and long-slit spectroscopy technique (e.g. Møller 2000; Møller et al. 2000; Weidinger et al. 2005), but here we take advantage of the integral field spectroscopy, which provides us images of the quasar and its surrounding at different wavelengths, including spectral regions where no extended line emission is expected, thus allowing for a careful PSF modelling. To subtract the PSF contribution we adopt a purely empirical method successfully tested in Borisova et al. (2016), using CUBEPSF, part of the CUBEXTRACTOR package (Cantalupo, in preparation). The algorithm consists in producing a pseudo-NB image with a spectral width of 150 spectral pixels ( $\sim 187 \text{ \AA}$ ) for each wavelength layer. In this way we construct empirical PSF images to be rescaled and subtracted at each wavelength layer. We rescale the flux in each empirical PSF image assuming that the quasar dominates the flux within a central region of  $5 \times 5$  pixels (corresponding to  $1 \times 1 \text{ arcsec}^2$ ). The rescaling factor between the flux in each layer and the empirical PSF images is computed using an averaged-sigma-clip algorithm in the central  $5 \times 5$  pixels region, to avoid any possible contamination by cosmic rays or other sources of noise in the individual pixels. We then cut from the rescaled empirical PSF images a central circular region with a radius of  $\sim 5$  times the seeing. As a last step, we mask any pixel with negative flux from this circular cut-out region, and we subtract it from the corresponding layer in the data cube. The spectral width of  $\sim 187 \text{ \AA}$  is a good compromise between minimizing the PSF variations within the waveband and maximizing the signal-to-noise ratio (SNR) in the empirical PSF images. We iterate the PSF removal procedure, masking the spectral interval associated with the nebulae to avoid oversubtraction. We also remove and mask continuum sources around the quasar, to be not affected by their residuals.

We note that, by construction, our PSF subtraction has uncertainties in the  $1 \times 1 \text{ arcsec}^2$  region used for the PSF rescaling, in particular due to the poor pixel sampling in the central region, where the nuclear emission has a very steep gradient. For this reason, any data on such scales are considered not reliable in the remaining of this work. We also note that our determination of the PSF may be potentially contaminated by the emission from the host galaxy in the central region used for the PSF modelling (see e.g. Fynbo, Møller & Warren 1999; Fynbo, Burud & Møller 2000 and Zafar et al. 2011 for a discussion). Indeed, the continuum emission from the host galaxy may affect the rescaling factor and cause systematic effects in the PSF subtraction. However, in our case, the high luminosity of the quasar ( $\lambda L_{1450\text{\AA}} \sim 3 \times 10^{46} \text{ erg s}^{-1}$ ; AB magnitude in the  $i$  band,  $i\text{-mag} = 19.75^2$ ) in comparison to the expected luminosity of its host-galaxy at  $z \sim 5$ , motivate our assumption that the quasar dominates the flux budget within the  $1 \times 1 \text{ arcsec}^2$  used for the PSF rescaling. In addition, as discussed in Appendix A, within the deep C IV trough, we do not see any evidence for a prominent emission

from the host galaxy, further supporting that the host galaxy does not contribute significantly to the observed emission.

As shown in the next section, our empirical PSF subtraction method produces good results, unveiling low-SB flux around the quasar on large scales.

After quasar PSF subtraction has been performed, we remove any other possible continuum sources for each spaxel in the cube using a fast median-filtering approach, following Borisova et al. (2016).

### 2.2.2 3D detection and extraction

We use the PSF-subtracted and continuum-subtracted cube for the detection and the extraction of extended line emissions. This task has been performed using CUBEXTRACTOR (Cantalupo, in preparation), a 3D automatic extraction software based on a 3D extension of the connected-labelling-component algorithm with union finding of classical binary image analysis (e.g. Shapiro & Stockman 2001).

We first extracted a sub-cube ( $\Delta\lambda \sim 1000 \text{ \AA}$ ) from the processed data cube with wavelength range containing the expected Ly $\alpha$  line and, before searching for detection, we apply a spatial Gaussian filtering of  $0.5 \text{ arcsec}$  (without smoothing in wavelength) to bring out extended but narrow features. An object is detected if it contains a number of connected voxels (individual spatial and spectral elements in MUSE cube) above a user-defined threshold (we used a threshold of 10 000 connected voxels) above a minimum SNR of 2.5, after data cube being smoothed. The 3D-segmentation-mask produced by CUBEXTRACTOR is then used for photometry, optimally extracted images, velocity, and velocity dispersion maps.

## 3 RESULTS

Following the data processing steps described in Section 2 we detect a Ly $\alpha$  nebula, extended on circumgalactic scales around the quasar J1605–0112, at  $z \sim 4.92$ .

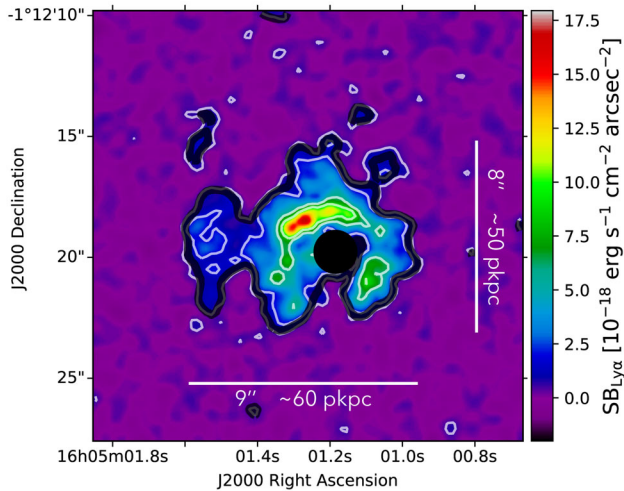
The nebula has been detected with a threshold of 10 000 connected voxels, and the resulting ‘optimally extracted image’, obtained using the 3D-segmentation-mask produced by CUBEXTRACTOR, is shown in Fig. 1. The black thick contour indicates the projection on the sky-plane of the 3D-mask, corresponding to an SB of  $\sim 6 \times 10^{-19} \text{ erg s}^{-1} \text{ cm}^{-2} \text{ arcsec}^{-2}$ . A single spectral layer, corresponding to the central wavelength of the Ly $\alpha$ , has been added – by means of the union operator – to the 3D-mask for display purposes.

The maximum spectral width defined in the 3D-mask corresponds to  $\Delta\lambda = 38.75 \text{ \AA}$ . The total Ly $\alpha$  flux, obtained integrating the 3D-mask along the wavelength and the spatial directions, is  $F_{\text{Ly}\alpha} \sim 1.7 \times 10^{-16} \text{ erg s}^{-1} \text{ cm}^2$ , corresponding to a luminosity  $L_{\text{Ly}\alpha} \sim 4.4 \times 10^{43} \text{ erg s}^{-1}$ .

The Ly $\alpha$  emission projected on the sky-plane appears to extend over about  $9 \text{ arcsec}$  in RA and  $8 \text{ arcsec}$  in Dec., i.e. respectively  $60$  and  $50 \text{ pkpc}$  at this redshift. The position of the quasar is marked by the black dot, whose dimension indicates the region used to normalize the PSF for each layer, i.e. the spatial region where the flux is dominated by the quasar (see Section 2.2.1). This image has been obtained selecting all the voxels detected by the 3D-segmentation-mask and integrating their fluxes along the wavelength direction.

Operationally speaking, this approach is equivalent to obtain various pseudo-NB images where the filter spectral size is optimized for each spaxel, to maximize the SNR of the object. However, noise estimation for an image obtained through this method is critical, as the noise associated with each spaxel depends on the number of layers that define it (i.e. the spectral size of pseudo-NB filter). This

<sup>2</sup> Computed from our MUSE data cube with an aperture of  $3 \text{ arcsec}$  in diameter, assuming an SDSS  $i$ -band filter. This value is not corrected for intrinsic absorption.



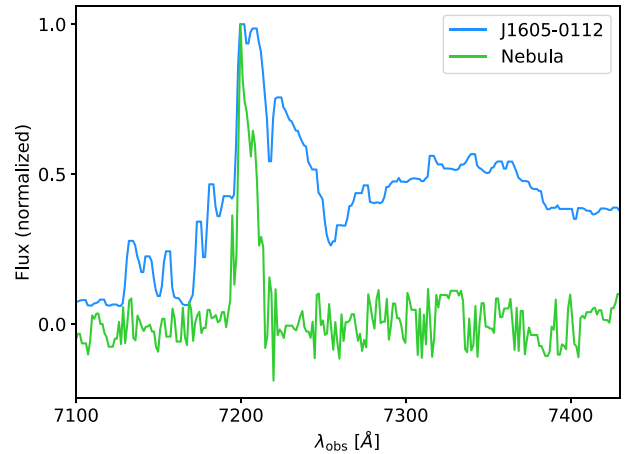
**Figure 1.** ‘Optimally-extracted’ image of the Ly $\alpha$  emission around J1605–0112. This image has been produced applying a 3D-mask (see discussion in Section 2.2) to the PSF and continuum subtracted MUSE data cube. The position of the quasar is marked by a black dot. For display purposes, we have added – by means of the union operator – to the object 3D-mask one wavelength layer of the cube corresponding to the central wavelength of the nebulae. The spatial projection of the 3D-mask is indicated by the black thick contour that correspond to an SB of  $\sim 5 \times 10^{-19} \text{ erg s}^{-1} \text{ cm}^{-2} \text{ arcsec}^{-2}$ . The thin contours indicate the propagated SNR in the images, with contour levels representing SNR = 2, 5, 10, 15. The white horizontal (vertical) bar indicates a physical scale of  $\sim 60$  pkpc ( $\sim 50$  pkpc). Both horizontal and vertical bars are representative of the maximum linear projected size of the nebula on the RA and Dec. axis.

results in the optimally extracted flux map being affected by a noise that depends on the spatial position. To address this issue, following Borisova et al. (2016), we estimated the noise and the SNR for each pixel in the map by propagating the variance and taking into account the number of layers contributing to each pixel. Using propagated variances and the integrated flux we then created the SNR contours represented by the white thin contours in Fig. 1.

In Fig. 2 we compare the Ly $\alpha$  profile of J1605–0112 with the Ly $\alpha$  profile of a region around the peak of the emission of the nebula (within an aperture of  $\sim 1$  arcsec). The Ly $\alpha$  emitted by the nebula show a much narrower profile (FWHM  $\sim 1000 \text{ km s}^{-1}$ ) than the broad and complex Ly $\alpha$  profile coming from the BAL quasar. Given such a difference between the widths and shapes of the spectra, we are confident that the Ly $\alpha$  emission from the nebula emission in close proximity to the quasar cannot constitute a PSF subtraction residual.

It is worth stressing here that the ‘optimally extracted image’ of the Ly $\alpha$  emission shown in Fig. 1 is different from a common standard NB image, because of the different number of spectral layers composing the pixels within the map. This type of approach, enabled by the IFU technique, allows us to reveal both kinematically narrow and broad features that would have been either lost in the noise or underestimated in an NB image with a single, fixed width.

We also produced a ‘plain’ pseudo-NB image (i.e. by averaging a fixed number of spectral channels, without any tailoring) of the Ly $\alpha$  emission. To this aim, we average over  $38.75 \text{ \AA}$  around the  $\lambda_{\text{Ly}\alpha}$  (rest-frame), that is the maximum spectral width of the nebula as defined by the 3D-mask (see Table 2). Such a pseudo-NB image recovers  $\sim 80$  per cent of the total flux revealed by the 3D-extraction, resulting in an apparent slightly smaller projected linear size of the nebula (about  $6.5 \text{ arcsec}$  in RA and  $7 \text{ arcsec}$  in Dec., respectively



**Figure 2.** Comparison of the Ly $\alpha$  line shapes between the BAL quasar J1605–0112 (blue line) and the peak of the detected Ly $\alpha$  nebula (green line). The spectrum of the nebula has been extracted within a circular aperture of  $\sim 1$  arcsec radius centred around the peak of its emission from the PSF and continuum subtracted final data cube. Both spectra have been normalized to their maximum to facilitate the comparison. The Ly $\alpha$  emission of the nebula appears to be much narrower than the broad Ly $\alpha$  of the quasar J1605–0112, ensuring that the peak of the nebula close to the quasar position is not due to PSF residuals.

**Table 2.** Properties of the Ly $\alpha$  nebula around J1605–0112.

Size <sup>a</sup> (kpc)	$\Delta\lambda^b$ ( $\text{\AA}$ )	Luminosity ( $\text{erg s}^{-1}$ )	FWHM <sup>c</sup> ( $\text{km s}^{-1}$ )	$C_{\text{IV}} / \text{Ly } \alpha^d$ ( $2\sigma$ )
50–60	38.75	$4.4 \times 10^{43}$	900	$< 0.06$

<sup>a</sup>Maximum linear projected size obtained from the 3D-mask.

<sup>b</sup>Maximum spectral width obtained from the 3D-mask, and used for the pseudo-NB image.

<sup>c</sup>Spatially averaged FWHM.

<sup>d</sup>Limit on the  $C_{\text{IV}} / \text{Ly } \alpha$  line ratio ( $2\sigma$ ).

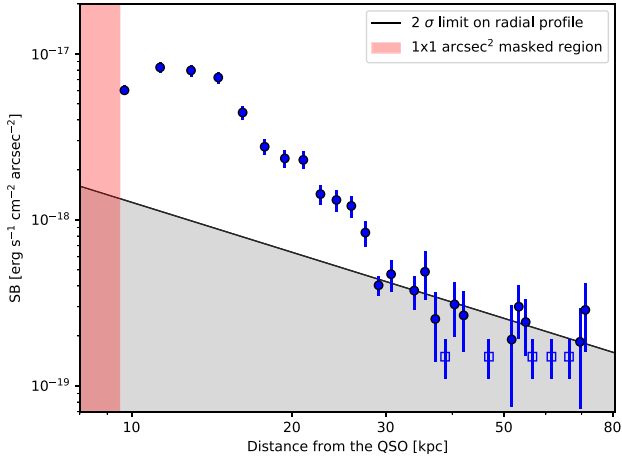
42 and 45 pkpc at this redshift; see Appendix C). For a consistent comparison with previous works, and to accurately characterize the noise and SB limits, we used the pseudo-NB image instead of the 3D optimally extracted image to extract the circularly averaged SB profile of the Ly $\alpha$  nebula, described in the next paragraph.

### 3.1 Surface brightness profile

We computed the circularly averaged SB profile of the Ly $\alpha$  nebula surrounding J1605–0112, at  $z = 4.92$ , by using the pseudo-NB image of Ly $\alpha$ , extracted as discussed above. The resulting profile, as a function of projected physical distance from the quasar, is presented in Fig. 3(a) (blue dots), where the expected  $2\sigma$  Gaussian SB limit for each aperture is also shown (grey shaded area).

To avoid the cosmological SB dimming effect and facilitate comparisons with similar objects at different redshifts, we computed the circularly averaged SB profile in units of  $\text{SB} \times (1+z)^4$ , i.e. cosmologically invariant brightnesses. The resulting profile is shown in Fig. 4(a), where the redshift of the quasar,  $z = 4.92$ , has been used as scaling value. Along with our results we also show the averaged profiles of two samples of quasars at redshift  $z \sim 3\text{--}4$  (the mean redshift is  $z \sim 3.2$  in both samples) by Borisova et al. (2016) and Arrigoni-Battaia et al. (in preparation) (violet and red solid lines, respectively). Also these profiles are reported in cosmologically





**Figure 3.** The circularly averaged SB profile (blue filled dots) of the Ly $\alpha$  nebula is shown as a function of the projected physical distance from J1605–0112. Open blue squares indicate negative values. Error bars represent  $1\sigma$  errors on the mean SB level measured in a given annulus. The SB profile has been extracted using a pseudo-NB image centred at the wavelength corresponding to the Ly $\alpha$  of the quasar, with a fixed width of 38.75 Å (see discussion in Section 3.1). The background flux level of the pseudo-NB image is consistent with zero. The grey shaded area is an estimate of the  $2\sigma$  Gaussian noise associated with the SB profiles. The pink shaded area is indicative of the  $1 \times 1$  arcsec<sup>2</sup> region used to rescale the empirical PSF models (see Section 2.2.1).

invariant SB units, using the mean redshift of the samples as scaling value.

Both Figs 3 and 4(a) show that the circularly averaged SB radial profile of the nebula surrounding J1605–0112 extends up to  $\sim 30$  pkpc from the quasar; the mean flux measured in annuli at larger radii is consistent with the noise ( $2\sigma$ ) limit. However, Fig. 4(a) shows that at distances consistent with the maximum extension of the radial profile of our nebula ( $\sim 30$  pkpc), the sensitivity provided by our deep MUSE observations would be enough to detect with

high SNR ( $\gtrsim 5$ ) any more extended emission, such as that expected by the averaged radial profiles of similar quasars at lower redshift.

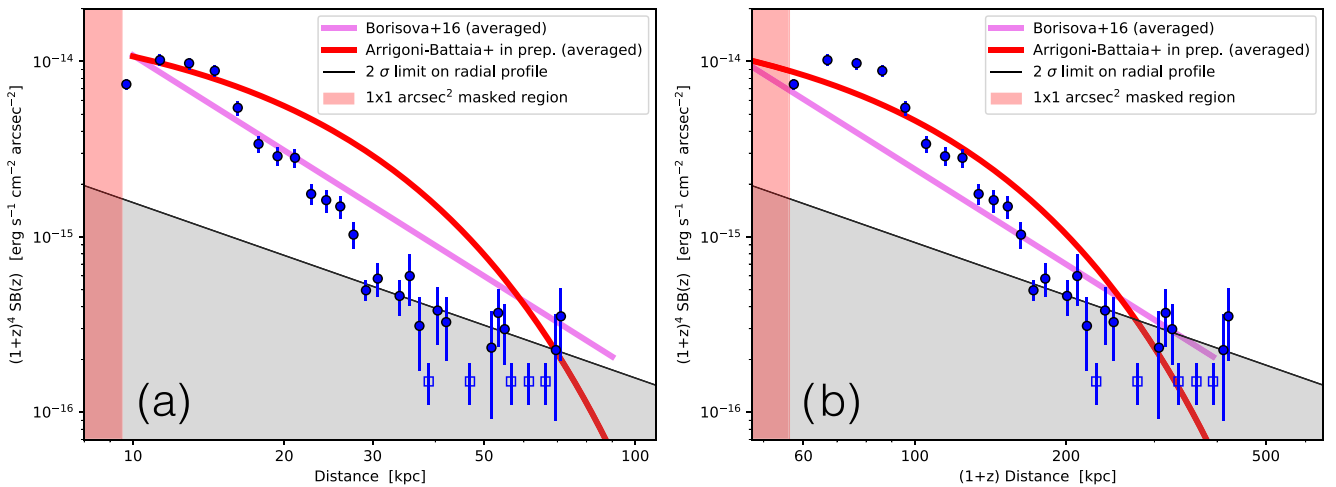
This suggests that the Ly $\alpha$  nebula detected in the surrounding of J1605–0112, at  $z \sim 5$  is intrinsically smaller than typical Ly $\alpha$  nebulae observed with MUSE around luminous quasars at lower redshifts,  $z \sim 3$ –4. We stress that such size discrepancy is intrinsic and not an artefact of the cosmological SB dimming; indeed the radial profiles have been expressed in invariant units, rescaling the quantities with redshift as discussed above.

In Fig. 4(b) we report the same SB radial profiles of Fig. 4(a), but in this case we show them as a function of the projected comoving distance from the quasar. Using comoving distances, instead of physical, we can account for the cosmological growth of DM haloes, whose virial radii, for a fixed halo mass, scale as  $\propto (1+z)^{-1}$  (Barkana & Loeb 2001). Fig. 4(b) shows that the intrinsic size discrepancy between our nebula at  $z \sim 5$  and typical nebulae at  $z \sim 3$ –4 is greatly reduced, or even disappears, when referring to comoving distances from the quasar. This suggests that our nebula appears to be less extended because haloes hosting quasars at high redshift are, on average, intrinsically smaller. In Section 4 we analyse possible origins and consequences of such finding.

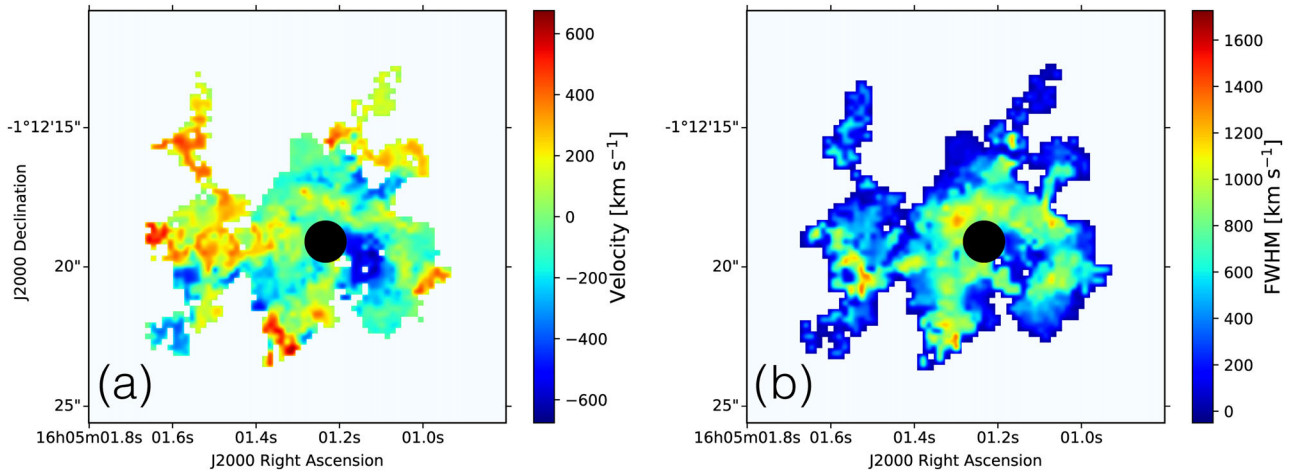
### 3.2 Kinematics

We used the 3D-segmentation-mask, defined in Section 2.2, to also produce two-dimensional maps of the first and second moments of the flux distribution in its spectral domain. Such maps give indication of, respectively, the centroid velocity and the width of the emission line (the velocity dispersion) at each spatial location.

In Fig. 5(a) we report the map of the first moment of the Ly $\alpha$  flux distribution, i.e. the shift of the velocity centroid with respect to the peak of the integrated emission of the nebula. The velocity map of Ly $\alpha$  emission does not show any clear evidence of rotation nor any coherent kinematic patterns, in line with similar recent MUSE observations at lower redshift (Borisova et al. 2016). However we note that, due to its resonant nature, Ly $\alpha$  is usually not a good tracer



**Figure 4.** (a) The radial profile shown in Fig. 3 here is expressed in units of  $SB \times (1+z)^4$  to avoid the cosmological SB dimming effect. Redshift-corrected averaged profiles of two samples of quasars observed with MUSE at mean redshift  $z \sim 3.2$  (Borisova et al. 2016; Arrigoni-Battaia et al., in preparation) are also shown for comparison. With a maximum linear projected size of  $\sim 30$  pkpc from the quasar, our nebula appears to be intrinsically smaller than typical nebulae observed around similar quasars at lower redshift. (b) The same radial profile is shown as a function of the projected comoving distance from the quasar, to take into account the cosmological growth of dark matter haloes. In this case, the intrinsic size discrepancy revealed in (a) is reduced (see Sections 3.1 and 4 for an interpretation of such findings).



**Figure 5.** (a) Velocity map traced by the  $\text{Ly}\alpha$  emission around J1605–0112. The black circle indicates the position of the quasar. The velocity map does not show any clear evidence of rotation or other coherent kinematic patterns, in line with similar recent MUSE observations at lower redshift (e.g. Borisova et al. 2016). (b) Velocity dispersion map showed in units of Gaussian-equivalent FWHM (i.e. 2.35 times the velocity dispersion) for consistency with previous works in the literature. The black circle indicates the position of the quasar. The  $\text{Ly}\alpha$  velocity dispersion map shows a particularly high broadening of the line ( $\text{FWHM} > 1000 \text{ km s}^{-1}$ ), especially in the inner regions of the CGM, at  $\gtrsim 10 \text{ pkpc}$  from the quasar.

of kinematics as any coherence on  $10 \text{ pkpc}$  scales may be disrupted by radiative transfer effects (Cantalupo et al. 2005).

In Fig. 5(b) we present the map of the second moment of the  $\text{Ly}\alpha$  flux distribution, i.e. the velocity dispersion (for consistency with previous works in the literature we show the Gaussian-equivalent FWHM derived by multiplying the second moment by 2.35). The  $\text{Ly}\alpha$  velocity dispersion map shows a particularly high broadening of the line ( $\text{FWHM} > 1000 \text{ km s}^{-1}$ ), especially in the inner regions of the CGM, at  $\sim 10 \text{ pkpc}$  from the quasar. Such a high velocity dispersions are in contrast with previous MUSE observations by Borisova et al. (2016), which found narrower nebulae with  $\text{FWHM} \sim 500\text{--}700 \text{ km s}^{-1}$  surrounding 16 (out of 17) radio-quiet quasars, and are more similar to what found in nebulae around radio-loud systems (Borisova et al. 2016). Two notable exceptions are the nebula #6, surrounding J0124+0044 at  $z \sim 3.8$  (Borisova et al. 2016), and the Enormous Lyman-Alpha Nebula surrounding J1020+1040 at  $z = 3.164$  (Arrigoni Battaia et al. 2018), both showing a broad  $\text{Ly}\alpha$  emission ( $\text{FWHM} \sim 1000 \text{ km s}^{-1}$ ) on similar scales ( $\sim 10 \text{ kpc}$  from the quasar) as for our nebula.

### 3.3 C IV/Ly $\alpha$ line ratio

After hydrogen  $\text{Ly}\alpha$  line,  $\text{C IV } 1549 \text{ \AA}$  and  $\text{He II } 1640 \text{ \AA}$  are the brightest UV lines found within quasar-powered nebulae, and their line ratios have been often used to constrain the ionization parameter, gas density, and metallicity, providing information on the origin and physical properties of the emitting gas (see Nagao, Maiolino & Marconi 2006; Arrigoni Battaia et al. 2015b, 2015a; Prescott et al. 2015, and references therein). For instance, giant  $\text{Ly}\alpha$  nebulae around high- $z$  radio galaxies typically show high  $\text{C IV/Ly}\alpha$  and  $\text{He II/Ly}\alpha$  (ratios  $\gtrsim 0.1$ ; see Villar-Martín et al. 2007), which are currently explained by high metallicity gas whose emission is driven by ionization and jet–ambient interaction. We searched only for possibly extended  $\text{C IV}$  within our MUSE data cube. We note that, at  $z = 4.92$ , the MUSE wavelength coverage does not enable us to study the  $\text{He II}$  emission line. In order to search for spatially extended  $\text{C IV}$  emission we used the  $\text{Ly}\alpha$ -defined 3D-mask produced by CUBEXTRACTOR. Following the analysis of Borisova et al. (2016) we scanned spectral layers around the location of  $\text{C IV}$ , by shifting

the mask along the wavelength direction of the PSF and continuum subtracted data cube. Using the same 3D-mask defined for the  $\text{Ly}\alpha$  avoids any aperture effects in the line ratios. As the  $\text{Ly}\alpha$  is expected to be intrinsically brighter and more extended than  $\text{C IV}$ , this approach gives us very conservative limits in case of non-detection. Considering the possibility of radiative transfer-induced velocity shift between the  $\text{Ly}\alpha$  line emission and the  $\text{C IV}$  we have scanned spectral layers within a window of  $3000 \text{ km s}^{-1}$  on both side of the expected  $\text{C IV}$  line. We then compute the total flux of the voxels associated with the 3D-mask at each spectral location.

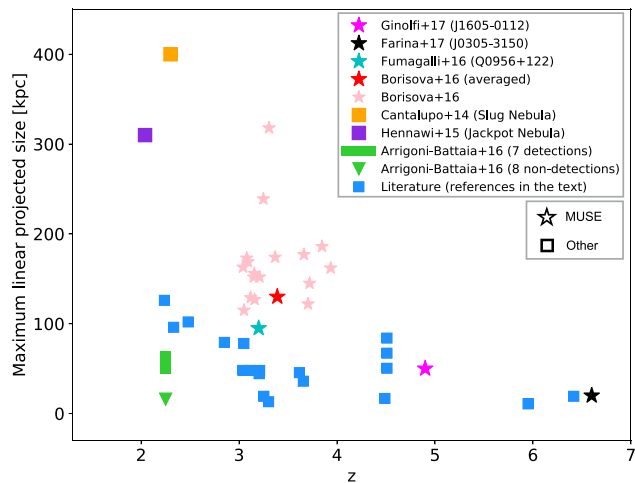
With this approach we do not detect any  $\text{C IV}$  extended emission around J1605–0112, within a  $1\sigma$  limit of  $\text{SB}_{\text{C IV}} (1\sigma) \sim 4.5 \times 10^{-18} \text{ erg s}^{-1} \text{ cm}^{-2}$ . This yields a  $2\sigma$  upper limit on the  $\text{C IV/Ly}\alpha$  ratio of  $\sim 0.06$ , consistently with what found in the majority of radio-quiet quasars at lower redshifts (e.g. Borisova et al. 2016). In the Appendix D, we show a  $30 \text{ \AA}$ -width NB image of the  $\text{C IV}$ , clearly showing that the  $\text{C IV}$  emission is undetected.

## 4 DISCUSSION

### 4.1 The nebula small size relative to lower redshift nebulae observed with muse

With a maximum projected linear size of  $\sim 60 \text{ physical kpc}$  (pkpc), the  $\text{Ly}\alpha$  nebula around J1605–0112, at  $z \sim 5$ , appears to be smaller than typical nebulae ubiquitously detected with MUSE around quasars at lower redshifts ( $z \sim 3\text{--}4$ ), extended on linear sizes larger than  $100 \text{ pkpc}$ , up to  $\sim 300 \text{ pkpc}$  (see Borisova et al. 2016).

As discussed in Section 3.1, we verified that such a size discrepancy is not a trivial effect of the SB cosmological dimming. Fig. 4(a) shows that, after correcting for the cosmological dimming effect [i.e. by expressing SB in units of  $\text{SB} \times (1+z)^4$ ], our  $\sim 5 \text{ h}$  on-source of MUSE integration provides a sensitivity level such that we should detect with high SNR ( $\gtrsim 5$ ) any typical emission at physical distances  $\gtrsim 30 \text{ pkpc}$  from the quasar, which would be expected by the averaged radial profiles of similar nebulae at lower redshift. In the following we provide two different scenarios whose combination is most likely the correct way to interpret such a physical size discrepancy.



**Figure 6.** The maximum projected sizes of a compilation of Ly $\alpha$  nebulae surrounding radio-quiet quasars is shown as a function of redshift. Star symbols refer to observations taken with MUSE, while squares to observations taken with NB filter technique (triangles refer to non-detections). Although the sample is affected by some caveats (see Section 4), this qualitative analysis suggests a positive trend between the sizes of Ly $\alpha$  nebulae and the cosmic time within the redshift range  $z \sim 2-7$ .

Cosmological theories of structure formation and, more in general the  $\Lambda$ CDM paradigm, predict that dark matter (DM) haloes collapsing at redshift  $z$  have a (physical) virial radius  $r_{\text{vir}} \propto M^{1/3}(1+z)^{-1}$  (e.g. Barkana & Loeb 2001). Thus, DM haloes at higher redshifts, if they remain in the same mass range, are smaller than DM haloes at lower redshifts. In Fig. 4(b) we showed the same SB dimming-corrected radial profiles shown in Fig. 4(a), but as a function of projected comoving distance from the quasars. Expressing distances in comoving units, defined as  $d_{\text{comov}} = d_{\text{phys}} \times (1+z)$ , compensates the redshift evolution of virial radii accounting for the cosmological growth of DM haloes across the cosmic time. Fig. 4(b) shows that, when referring to comoving distances, the intrinsic size discrepancy between our nebula at  $z \sim 5$  and typical nebulae at  $z \sim 3-4$  is greatly reduced. This suggests that our nebula appears to be less extended because quasar hosting DM haloes at high redshift are smaller than haloes at lower redshifts, assumed that they are hosted on average by haloes with comparable masses in the redshift range 3–5.

Such a result would suggest a physical correlation between the size of Ly $\alpha$  nebulae and the sizes of DM haloes/structures around quasars. This conjecture is supported by the qualitative analysis of other observations of Ly $\alpha$  nebulae in an even wider redshift range. In Fig. 6 we report the sizes (maximum linear projected sizes) of a compilation<sup>3</sup> of Ly $\alpha$  nebulae (including non-detections)

<sup>3</sup> We updated the collection reported in the PhD thesis of Arrigoni-Battaia (table 1.2; [https://www.imprs-hd.mpg.de/49473/thesis\\_Arrigoni.pdf](https://www.imprs-hd.mpg.de/49473/thesis_Arrigoni.pdf)). The updated list includes data from: Heckman et al. (1991a,b); Bremer et al. (1992); Roettgering et al. (1997); van Ojik et al. (1996); Lehnert & Becker (1998); Bergeron et al. (1999); Fynbo, Burud & Møller (2000); Bunker et al. (2003); Weidinger, Møller & Fynbo (2004); Weidinger et al. (2005); Christensen et al. (2006); Barrio et al. (2008); Courbin et al. (2008); Smith et al. (2009); Goto et al. (2009); Matsuda et al. (2011); Willott et al. (2011); Adelberger et al. (2005); North et al. (2012); Goto et al. (2012); Humphrey et al. (2013); Roche, Humphrey & Binette (2014); Cantalupo et al. (2014); Hennawi et al. (2015); Arrigoni Battaia et al. (2016); Borisova et al. (2016); Fumagalli et al. (2016); Farina et al. (2017).

surrounding radio-quiet quasars as a function of redshift, suggesting a positive trend between the sizes of Ly $\alpha$  nebulae and the cosmic time within the explored redshift range, i.e.  $z \sim 2-7$ . This qualitative trend has to be taken with caution due to the following caveats: (i) maximum projected sizes have been measured through different methods and with different SB limits, (ii) many of the nebulae are asymmetrical and it is not obvious how to characterize them with a single physical size, (iii) the sample is biased towards detections, as many non-detections have not been published. We leave a more quantitative analysis of this trend, including a re-processing of data in the literature and a self-consistent calculation of the maximum projected sizes (using a common isophote reference) for a future work. Fig. 6 also shows an interesting increase in the scatter across the cosmic time, possibly due to a lack of observations for high-redshift quasars or, more interestingly, to galaxy formation processes: quasar host haloes at low- $z$  have a larger variety of properties (e.g. masses, environment, amount of cold gas, UV source luminosities) than their higher  $z$  counterparts.

Another interpretation, expected to contribute to the observed size discrepancy, may be ascribed to the BAL-nature of the quasar J1605–0112. BAL quasars are characterized by redder continua with respect to non-BAL quasars (see Maiolino et al. 2004). In particular J1605–0112 requires a reddening  $E_{B-V} = 0.03$  of SMC extinction curve in order to match the observed near-IR continuum slope (Reichard et al. 2003; Maiolino et al. 2004). This high reddening is often related to a high column density of dust in the outflowing circumnuclear medium around the active galactic nucleus (AGN), able to efficiently absorb the AGN-powered UV photons, able to suppress the escape fraction of the ionizing radiation necessary to induce fluorescence in the neutral circumgalactic gas.

#### 4.2 The origin of the high Ly $\alpha$ broadening

Although Ly $\alpha$ , due to its resonant nature, is not generally considered a good tracer of kinematics, the analysis of circumgalactic gas kinematics of the nebula surrounding J1605–0112, reported in Section 3.2, shows peculiar results. In particular the velocity dispersion map reported in Fig. 5(b) (where the Gaussian-equivalent FWHM of the line is shown at each spatial location) shows a particularly high broadening of the line ( $\text{FWHM} > 1000 \text{ km s}^{-1}$ ), especially in the inner regions of the CGM, at  $\sim 10 \text{ kpc}$  from the quasar, in agreement with the scales in Arrigoni Battaia et al. (2018). Such a high value of FWHM exceeds by a factor of  $\sim 2$  the typical FWHM observed in other works. For instance Borisova et al. (2016) reported observations of narrower nebulae, e.g. with  $\text{FWHM} \sim 500-700 \text{ km s}^{-1}$ , surrounding almost all the radio-quiet quasars in their sample (with the notable exception of the nebula #6, surrounding J0124+0044 at  $z \sim 3.8$ , showing a radial profile and a kinematics that looks similar to our nebula).

A possible interpretation of such high broadening of the Ly $\alpha$  line emission in the inner region of our nebula may be ascribed to the BAL nature of the quasar J1605–0112 (Maiolino et al. 2004). Indeed, the outflow inferred from the BAL features, may not be restricted to the circumnuclear region, but extend to large scales. For instance, based on a detailed spectral analysis of a few BAL quasars, Borguet et al. (2013) infer that in some cases the outflowing gas traced by the absorption troughs is distributed on scales of  $10-30 \text{ kpc}$ , hence revising the standard scenario in which BALs are only associated with outflows on small scales. Therefore the large FWHM of the Ly $\alpha$  may directly trace the large scale outflow associated with the BAL. Even if the Ly $\alpha$  is not directly tracing the outflow gas, as suggested by the null detection of the C IV line,

possibly pointing to a low gas metallicity, the outflow can anyhow introduce significant turbulence into the CGM which may result in the observed broadening.

## 5 CONCLUSIONS

Recent MUSE observations have revealed ubiquitous giant Ly  $\alpha$  nebulae around bright quasars at  $z \sim 3\text{--}4$ , extending up to  $\sim 300$  pkpc (e.g. Borisova et al. 2016; Fumagalli et al. 2016; Arrigoni Battaia et al. 2018). In this work we present deep ( $\sim 4$  h) MUSE observations of J1605–0112, a BAL quasar at higher redshift,  $z \sim 5$ . Our findings can be summarized as follows:

(i) we reveal a Ly  $\alpha$  nebula around J1605–0112. The Ly  $\alpha$  emission projected on the sky-plane appears to extend over  $\sim 8\text{--}9$  arcsec, i.e. 50–60 kpc at  $z = 4.92$  (Fig. 1);

(ii) after correcting for the cosmological SB dimming, we compared the circularly averaged SB radial profile of our nebula with the averaged radial profiles of other nebulae at lower redshift, observed with MUSE. The profile of our Ly  $\alpha$  nebula, at  $z \sim 5$ , shows a steeper slope than typical nebulae at lower redshifts and indicate a less extended gas distribution (Fig. 4a). It is important to note, however, that single nebulae at lower redshift show a large variability, also in terms of radial profiles;

(iii) the effect described in (ii) is greatly reduced when referring to comoving distances (Fig. 4b), which take into account the cosmological growth of DM haloes [whose virial radius is expected to evolve proportionally to  $1/(1+z)$ ]. This suggests that the discrepancy between the Ly  $\alpha$  extensions may be ascribed to the smaller size of DM haloes with comparable mass at higher  $z$ , implying an interesting empirical relation between the size of Ly  $\alpha$  nebulae and sizes of DM haloes around quasars. This conjecture is supported by a qualitative analysis of other observations of Ly  $\alpha$  nebulae in a larger redshift range (see Fig. 6);

(iv) while the velocity map does not show any clear pattern (Fig. 5 a; in line with similar observations at lower redshift), the velocity dispersion map shows a particularly high broadening of the line ( $\text{FWHM} > 1000 \text{ km s}^{-1}$ ) in the inner regions of the nebula (Fig. 5b).

We suggest that (iii) and (iv) findings may be favoured by the BAL nature of the observed quasar. Indeed BAL quasars are characterized by redder continua with respect to non-BALs (e.g. Reichard et al. 2003; Maiolino et al. 2004) and this is often associated with high dust column densities in the circumnuclear regions around the AGN. Such a large dust content may efficiently absorb the UV photons emitted by the AGN and suppress the escape fraction of the ionizing radiation responsible for producing the Ly  $\alpha$  fluorescence in the CGM, thus reducing the detectable size of Ly  $\alpha$  nebulae. Moreover, BALs have been observed to experience powerful outflows and in particular LoBALs (as J1605–0112) show extremely deep troughs of low-ionization lines (e.g. Maiolino et al. 2004), indicative of very large columns of outflowing gas. This is consistent with BALs tracing a phase of quasar feedback particularly effective in ejecting gas (e.g. Dunn et al. 2010). In this light, the broad Ly  $\alpha$  emission observed in the inner region (on  $\sim 10$  kpc scales) around our BAL quasar may trace both highly turbulent circumgalactic gas as well as outflowing material escaped from the galaxy.

## ACKNOWLEDGEMENTS

The authors would like to thank the anonymous referee for the useful comments. MG thanks F. G. Saturni for helpful discussions. The research leading to these results has received funding from the European Research Council (ERC) under the European Union's Seventh

Framework Programme (FP/2007-2013) / ERC Grant Agreement no. 306476. RM acknowledges support from the ERC Advanced Grant 695671 'QUENCH'. RM and S. Carniani acknowledge support from the Science and Technology Facilities Council (STFC). S. Cantalupo gratefully acknowledges support from Swiss National Science Foundation grant PP00P2\_163824. The work is based on observations collected at the European Organisation for Astronomical Research in the Southern hemisphere under ESO programme 095.B-0875(A). This research made use of *ASTROPY*, a community-developed core *PYTHON* package for Astronomy (Astropy Collaboration et al. 2013).

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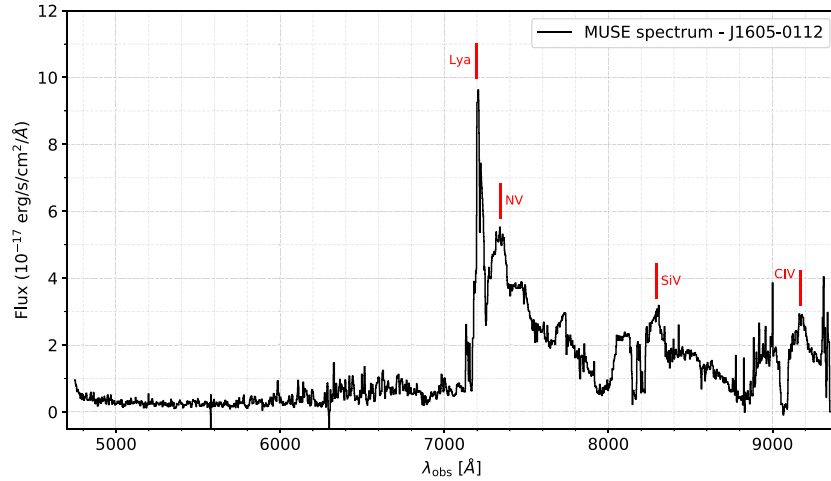


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## APPENDIX A: MUSE SPECTRUM OF J1605–0112 AND ITS BAL FEATURES

In Fig. A1 we show the 1D spectrum of the BAL quasar J1605–0112, at  $z = 4.92$ . The spectrum has been extracted in a 1.5 arcsec radius circular aperture, containing most of the flux, consistently with SDSS.

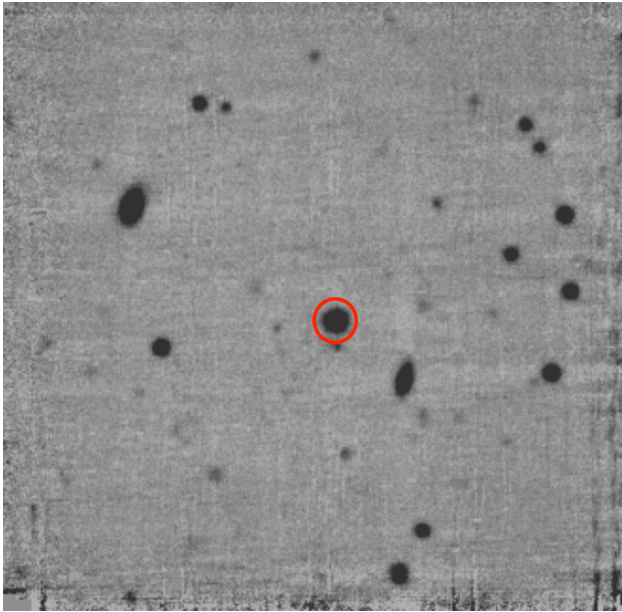
As visible from the spectrum, the deepest trough of the C IV, at  $\lambda_{\text{obs}} \sim 9070 \text{ \AA}$ , absorbs almost completely the light emitted by the quasar. In particular, using the spectral information of our MUSE observations, after computing a linear fit of the continuum emission in the wavelength range of the deep C IV trough, we estimate that the absorption leaves a residual flux which is only about 3 per cent of the continuum emission at those wavelengths. Such a residual light in the deepest absorption trough is consistent with a BAL covering fraction of the order of unity (e.g. Elvis 2000). It is unlikely that such 3 per cent residual emission is associated with the host galaxy as the radial profile is consistent with what observed in the quasar continuum outside the trough. Applying the same analysis to the other troughs of C IV, Si IV, and N V, we found less extreme residual fluxes, ranging from 20 per cent to 60 per cent. A more detailed analysis of the emission in the BAL troughs is beyond the scope of this paper and we leave it for future works.



**Figure A1.** One-dimensional spectrum of J1605–0112 obtained with MUSE. The spectrum has been extracted in a 1.5 arcsec radius circular aperture, containing most of the flux. The red lines indicate the location of UV emission lines, e.g. Ly  $\alpha$ , N v, Si v, C iv.

## APPENDIX B: WHITE-LIGHT IMAGE

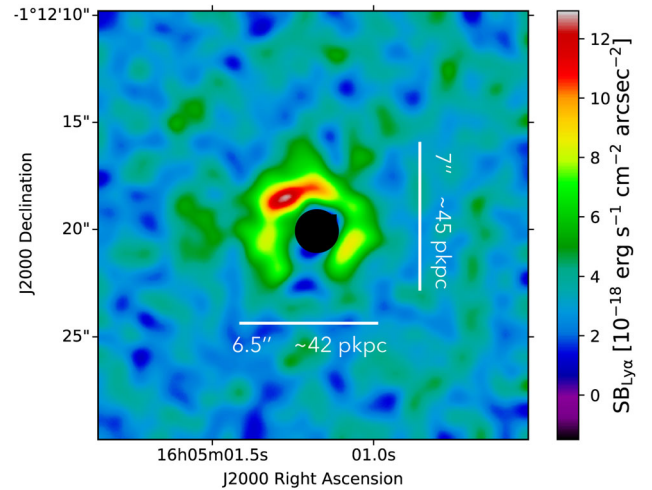
In Fig. B1 we show the whole  $1 \times 1$  arcmin<sup>2</sup> FOV of the final data cube as a ‘white-light’ image, obtained by collapsing the data cube along the wavelength direction.



**Figure B1.** White-light image of the whole  $1 \times 1$  arcmin<sup>2</sup> FOV of the final data cube. The red circle indicates the position of the quasar.

## APPENDIX C: PSEUDO-NB IMAGE OF THE LY $\alpha$ EMISSION

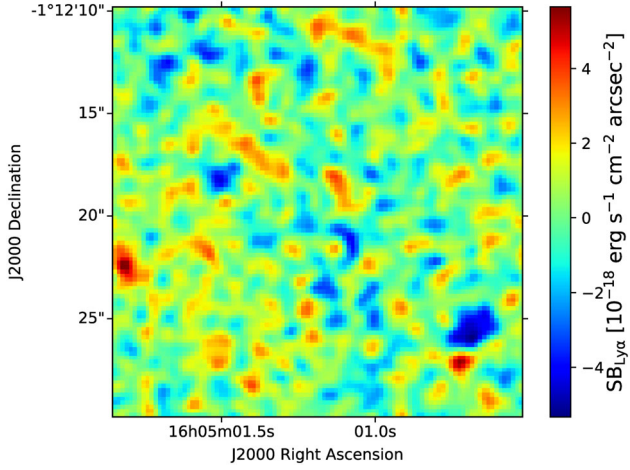
In Fig. C1 we show a pseudo-NB image of the Ly  $\alpha$  emission, obtained by averaging 38.75 Å (i.e. the maximum spectral width of the nebula as defined by the 3D-mask) around the rest-frame wavelength of Ly  $\alpha$  (see Table 2).



**Figure C1.** Pseudo-NB image of the Ly  $\alpha$  emission around J1605–0112, centred at the wavelength corresponding to the Ly  $\alpha$  of the quasar, with a width coincident with the maximum spectral width of the nebula as defined by the 3D-mask, i.e. 38.75 Å. To facilitate the comparison, the data are shown on the same scale of Fig. 1.

#### APPENDIX D: NON-DETECTION OF $\text{C IV}$ EMISSION

In Fig. D1 we show a pseudo-NB image of the  $\text{C IV}$  emission, obtained by averaging  $30 \text{ \AA}$  around the  $\text{C IV}$  central wavelength (rest-frame  $9170 \text{ \AA}$ ). It is evident that the  $\text{C IV}$  emission is undetected.



**Figure D1.** NB image of the  $\text{C IV}$  emission around J1605–0112, obtained by averaging  $30 \text{ \AA}$  around the  $\text{C IV}$  central wavelength (rest-frame  $9170 \text{ \AA}$ ). The  $\text{C IV}$  emission is clearly undetected.

This paper has been typeset from a  $\text{T}_{\text{E}}\text{X}/\text{L}_{\text{A}}\text{T}_{\text{E}}\text{X}$  file prepared by the author.